

Structure and Stability of X-Ray Irradiated Accretion Disk

¹Bari Maqbool, ²Ranjeev Misra, ³Naseer Iqbal, and ⁴Gulab Dewangan

**1, 3 Department of Physics, University of Kashmir, Srinagar, India.
2,4 Inter-University Centre for Astronomy and Astrophysics, Pune, India**

Abstract:

We present here the mathematical approach in calculating the structural changes which take place in the outer regions of the accretion disk due to X-ray irradiation.

It is shown here that an X-ray source powered by accretion, modifies the outer disc structure.

Our calculations for the transition radius and Circularization Radius in case of various low mass X-ray binaries show that the X-ray irradiation becomes dominant after transition radius only in some binary systems.

Structure of the X-ray Irradiated Accretion Disk:

The general Radiative Transport equation is given by

$$\frac{acT^4}{\tau} \approx \frac{acT^4}{\kappa\Sigma} \approx F(r), \quad \tau(r) > 1$$

where $F(r)$ is the flux of the heat generated internally by viscosity,

τ represents optical depth,

κ represents opacity,

Σ is the Surface Density.

or, The above equation can be written as

$$\sigma T^4 = F_0(r) [\tau]$$

Now taking X-ray Irradiation into consideration, the above equation modifies to (Dubus et.al. 1999)

$$\sigma T^4 = F_0(r) [\tau] + F_{ir}(r, h)$$

Where $F_{ir}(r,h)$ measures the flux due to X-ray Irradiation and is given by (Vrtelik et.al 1989)

$$F_{ir}(r, h) = \frac{f_2 L_x}{4\pi r} \frac{\partial h/r}{\partial r}$$

Taking into account both the opacities, we can write as:

$$\sigma T^4 = F_0(r) \left[\Sigma (\bar{\kappa}_{ff} + \bar{\kappa}_{es}) \right] + F_{ir}(r, h)$$

Solving the above equation with other Standard Model Equations, we get a differential equation given as:

$$\frac{\partial y}{\partial x} = \frac{y}{x} + A_p x^{-10} y^8 - B_p x^{\frac{25}{2}} y^{-12} - C_p x^{\frac{1}{2}} y^{-2}$$

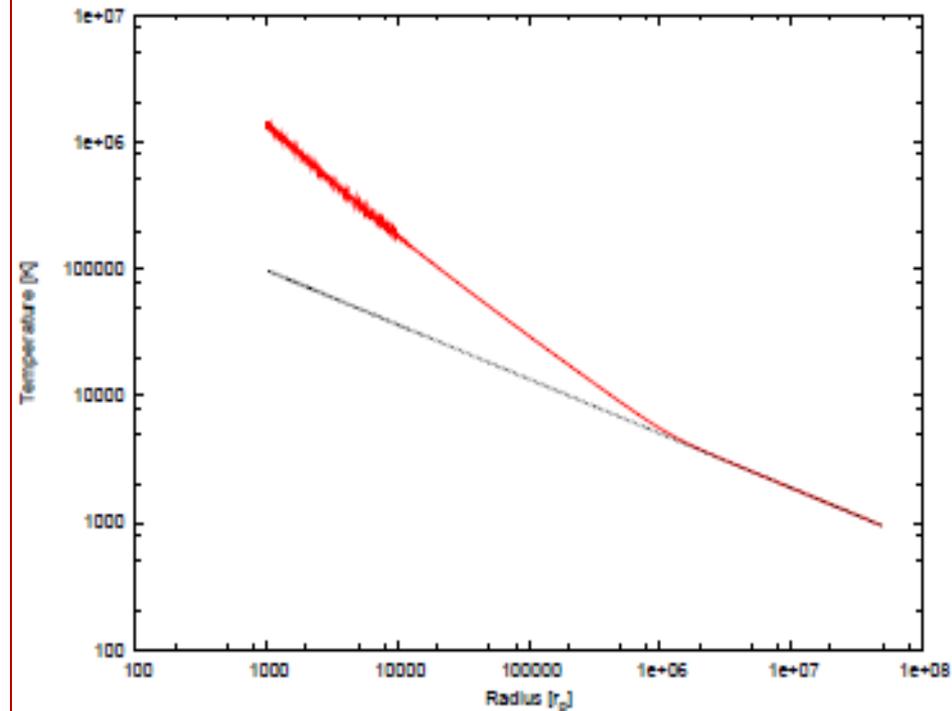
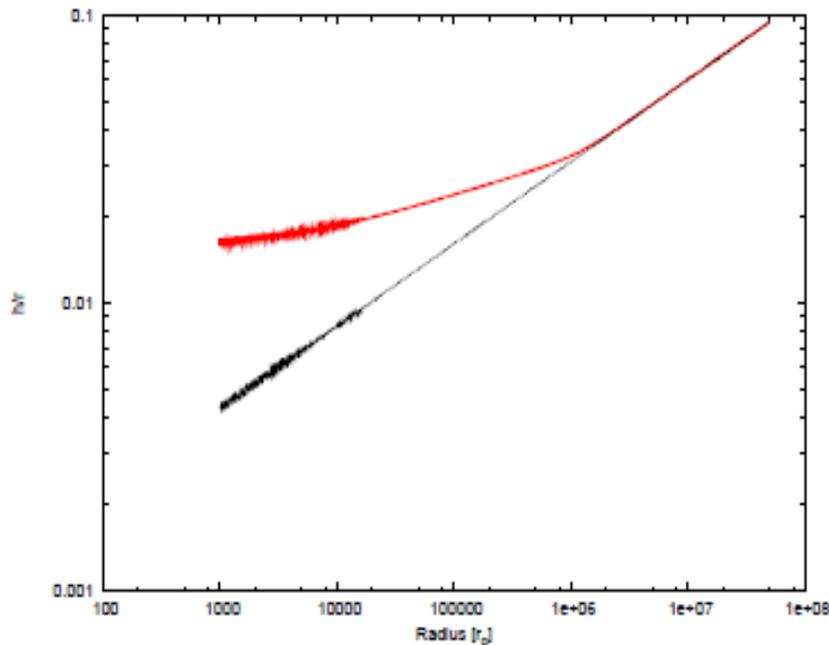
where

$$x = \frac{r}{r_g} \quad y = \frac{h}{r_g}$$

$$A_p = \frac{4\pi\sigma r_g^2 c^6}{f_2\eta\dot{M}_{in}} \left(\frac{\mu m_H}{k f_1}\right)^4 \quad B_p = \frac{12 \times 0.64 \times 10^{23}}{128} \left(\frac{\dot{M}_{out}^3}{\dot{M}_{in} r_g^3 c^9}\right) \left(\frac{1}{\pi^2 \alpha^2 f_2 \eta}\right) \left(\frac{k f_1}{\mu m_H}\right)^{\frac{7}{2}}$$

$$C_p = \frac{12}{32} \left(\frac{\dot{M}_{ou}^2}{\dot{M}_{in} r_g c}\right) \left(\frac{\bar{\kappa}_{es}}{\pi \alpha f_2 \eta}\right)$$

Solving equation (6) for fixed values of M , \dot{M} , α , and η , We get the plots as shown



Circularization Radius and Transition Radius

As is clear from the figure 2, that there is a sharp transition as, we go from the X-ray Irradiated region to the region where X-Ray Irradiation is not important. We have developed an expression for calculating this transition point, called as '*Transition Radius*, given by

$$r_{tr} = \left[\frac{(4\alpha)^{26/45}}{\pi^{2/45}} \right] \left(\frac{12 \times 0.64 \times 10^{23}}{128\alpha^2} \right)^{14/45} \left(\frac{7}{2f_2\eta} \right)^{8/9} (\dot{M})^{2/45} \left(\frac{m_H}{k} \right)^{11/9} c^{2/3} r_g^{11/9} \left(\frac{\mu}{f_1} \right)^{104/45}$$

$$r_{tr} = 1.4815 \times 10^{11} \times \left(\frac{\dot{M}}{10^{17}} \right)^{2/45} \left(\frac{r_g}{1.5 \times 10^5} \right)^{11/9} \left(\frac{\alpha}{0.1} \right)^{-28/45}$$

Table 1. Comparison between the Circularization radius, Tidal Radius and Transition Radius in case of Low Mass X Ray Binaries.

System	Type	P_{orb} (d)	$M_{1\odot}$	R_c ($10^{11} cm$)	R_T ($10^{11} cm$)	R_{tr} ($10^{11} cm$)	T (K)
3A1516-569*	NS	16.60	1.4	6.5	13.0	2.2	4900
1E1740-7-2942*	BHC	12.73	10	10	20	27.38	1400
GRS-1758-258*	BHC	18.45	10	14	28	27.38	1400
401811-17*	NS	24.0667	1.4	8.3	16.6	2.2	4900
GRS 1915 +105*	BHC	33.500	14	59	118	41.3	1100
GS2023 + 338*	BHC	6.4750	12	7.1	14.2	34.2	1200
CygX-2**	NS	9.84	1.780	3.2	6.4	3.0	4200
V395 CAR**	NS	9.02	1.44	3.1	6.2	2.3	4900
XTE J2123-058**	NS	0.25	1.415	.25	0.5	2.3	4900
2A 1822-371**	NS	0.23	0.97	0.2	0.4	1.47	6300

Viscous Time scale

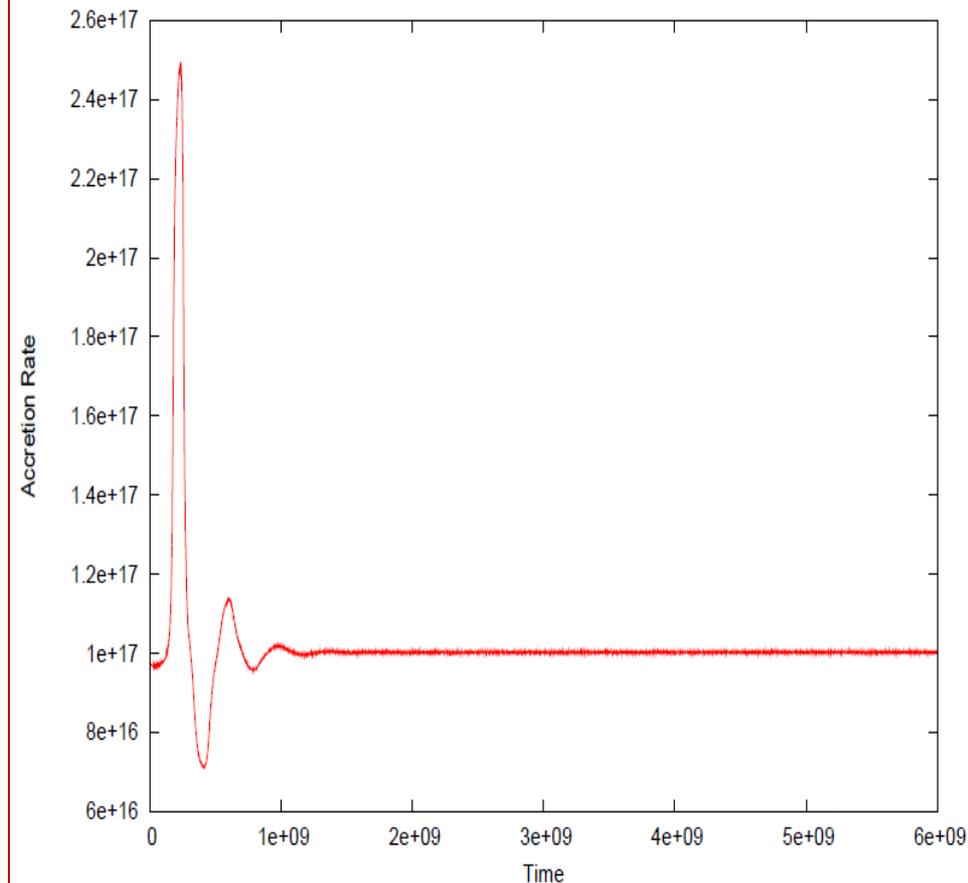
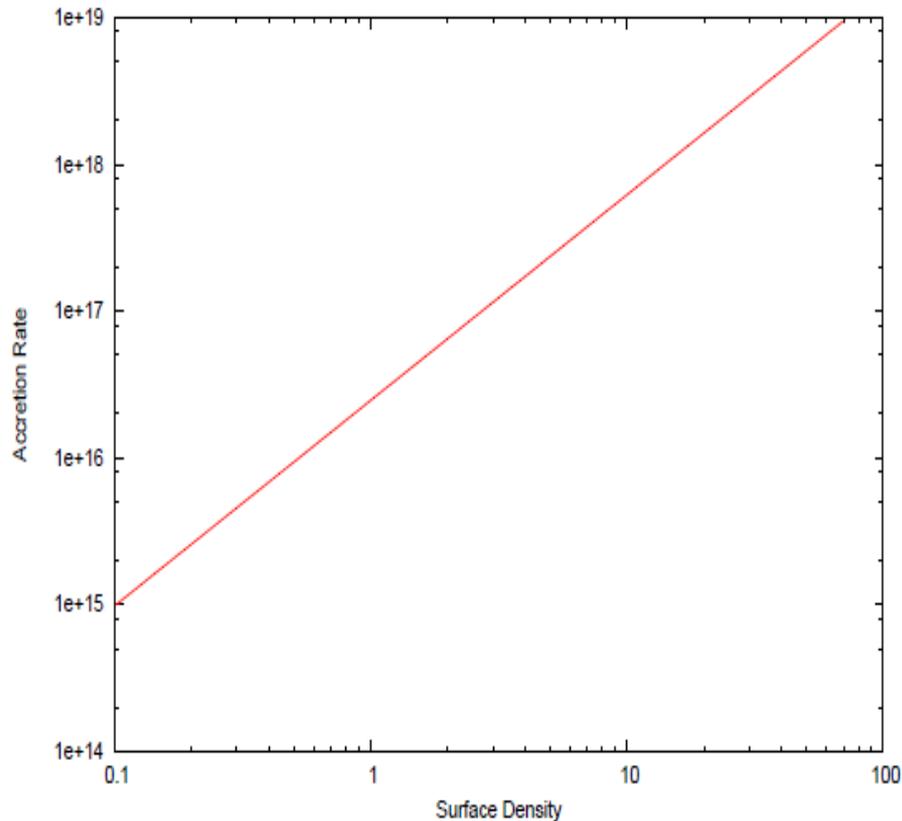
- It is the timescale on which matter diffuses through the disk under the effect of viscous torques and is given by:

$$t_{visc} = \frac{1}{2} \left(\frac{12 \times 0.64 \times 10^{23}}{128} \right)^{13/45} \left(\frac{7}{2\eta f_2} \right)^{10/9} \left[\frac{(4\sigma)^{37/45}}{(\alpha)^{71/45}} \right] \left(\frac{\dot{M}_{out}^{13/15}}{\dot{M}_{in}^{10/9}} \right) r_g^{16/9} c^{4/3} \left(\frac{\mu m_H}{k f_1} \right)^{41/18}$$

- For a Neutron star of mass 1.4 solar masses and uniform accretion rate of 10^{17} , the viscous timescale comes out to be 3.82×10^7 secs.

Stability:

For studying the stability of the such a disk, we plot a graph between the Accretion Rate and Surface density, Σ for a given value of radius and also considering That the inner accretion rate is same as the outer accretion rate. However , if the inner Accretion rate changes, the effect of the same on the outer Accretion rate on viscous timescale is depicted in the fig below.



References:

Frank, King, and Raine: *Accretion Power in Astrophysics* 3rd Edition,,
Cambridge. Univ. Press . 2002

Dubus G.,et.al, *Not .R Astron. Soc.* 303, 139-147 (1999)

Hoshi R. , Inoue H., *Astronomical Society Of Japan* 40, 421-434, (1988)

N.I.Shakura, R.A.Sunyaev , *Astronomy And Astrophysics*, 24, 337-355,(1973)

Vrtilek S.D., Raymond J.C, Garcia M.R., Verbunt F., Hasinger G., and

Kurster M., *Astronomy and Astrophysics*, 235, 162-173